

## NUMERICAL ANALYSIS OF DAMAGE DURING HOT FORMING

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### Abstract

The main aim of the presented research is the analysis of damage evolution in 16MnCrS5 steel during hot forming based on results obtained from finite element modelling. Particular attention is put on the interaction between dynamic recrystallization (DRX) and damage initiation at the matrix-inclusion interface. Moreover, a modified Gurson-Tvergaard-Needleman (GTN) model is proposed with the nucleation criterion taken from an extended Horstemeyer model, which predicts damage nucleation based on material softening due to the DRX and stress state in the material.

**Key words:** Damage, Dynamic recrystallization, Void nucleation, Numerical modelling

## 1. INTRODUCTION

Hot forming of metals is commonly used in the industry to produce semi-finished and finished products. It enables to lower the loads needed for the forming operation due to the reduction in yield stresses. Hot forming conditions also improve the material quality at the microscale because of the grain refinement due to recrystallization (Madej, et al., 2016). During hot forming, the ductility, toughness and material strength are improved (Muszka et al., 2014). Also, high temperature enables diffusion processes, which can remove/reduce chemical inhomogeneity (Pietrzyk & Kuziak, 2012; Milenin et al., 2017). In addition, the hot forming enables to reduce unavoidable casting defects, such as porosity (Saby et al., 2015). Unfortunately, the casting process also introduces non-metallic inclusions (impurities) into the material volume, which can cause damage (figure 1) during later hot forming, i.e. rolling (Malkiewicz & Rudnik, 1963; Ervasti & Stahlberg, 2005). Such inclusions, depending on their type, can

behave differently during hot forming. They can be rigid or deform with the matrix (Ervasti & Stahlberg, 2005).

Different investigations regarding the influence of non-metallic inclusions on material damage were performed in the past. Most of the literature focuses on experimental analysis of damage in cold forming (Requena et al., 2014; Aşık et al., 2019) and modelling of the observed phenomena (Requena et al., 2014; Liu et al., 2018; Tang et al., 2017; Santos et al., 2018). Only a limited amount of references focuses on damage in hot forming processes (Luo, 2001; Cheng et al., 2017).

Thus, the aim of this work is to analyse the evolution of damage during hot forming, including several phenomena that interact and influence the damage behaviour. Primarily, an interaction between the dynamic recrystallization and damage initiation and propagation at the matrix-inclusion-interface is numerically analysed in the paper.

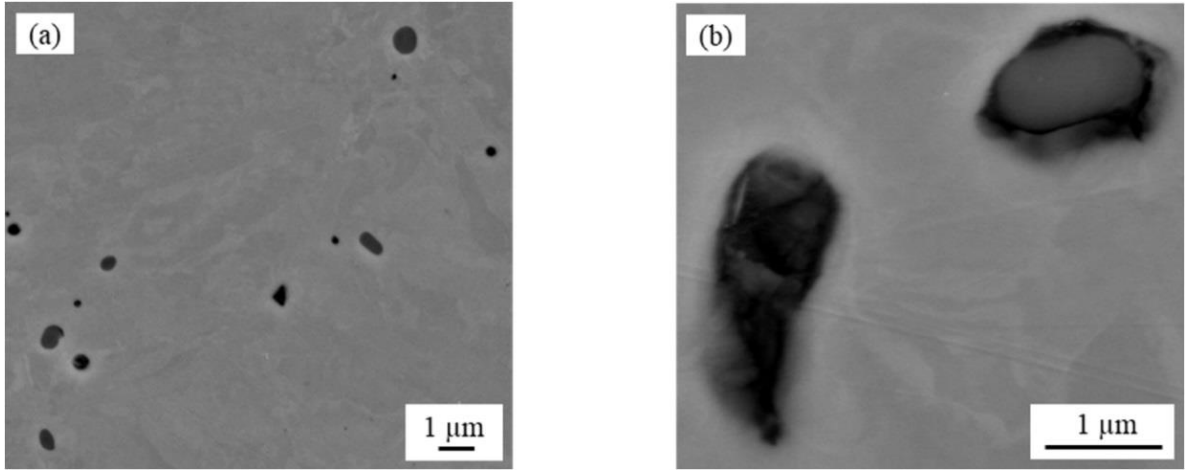


Fig. 1. Examples of the de-cohesion of steel matrix from the inclusion after material deformation.

## 2. FINITE ELEMENT MODEL FOR DAMAGE EVOLUTION

An axisymmetric numerical model, representing the material matrix with a non-metallic inclusion, was developed in Abaqus. The inclusion diameter is 4.25  $\mu\text{m}$  and covers the biggest inclusion size present in the investigated material. To avoid the influence of matrix size on the quality of obtained results, the matrix size is 10 times larger than the inclusion diameter. To simulate the damage evolution during deformation, temperature-dependent cohesive zone elements are located between the matrix and inclusion. The cohesive elements thickness is set to 0.01 $\mu\text{m}$ . The matrix and inclusion are discretized with 4-node bilinear axisymmetric elements (CAX4R) and the cohesive zone is meshed with 4-node axisymmetric cohesive elements (COHAX4). An illustration of the Representative Volume Element (RVE) is presented in figure 2.

In the current investigation, the inclusion is modelled as a rigid body, while the matrix has elastoplastic material properties. The matrix is assumed as an isotropic solid with a hardening model that considers dynamic recrystallization.

### 2.1 Dynamic recrystallization model

The average flow stress in the matrix is based on the model due to Beynon and Sellars equations, which simulate the softening of the matrix due to dynamic recrystallization during the deformation (Beynon & Sellars, 1992):

$$\sigma = \begin{cases} \sigma_0 + (\sigma_{ss0} - \sigma_0)[1 - \exp(-C \varepsilon)]^m, & \varepsilon < \varepsilon_{cr} \\ \sigma_0 + (\sigma_{ss0} - \sigma_0)[1 - \exp(-C \varepsilon)]^m - (\sigma_{ss0} - \sigma_{ss1})X, & \varepsilon \geq \varepsilon_{cr} \end{cases} \quad (1)$$

$$\sigma_0 = A_1 \sinh^{-1} \left[ \left( \frac{Z}{A_2} \right)^{A_3} \right] \quad (2)$$

$$X = 1 - \exp \left[ -k' \left( \frac{\varepsilon - \alpha \varepsilon_p}{\varepsilon_p} \right)^m \right] \quad (3)$$

$$\varepsilon_p = A_4 Z^q \quad (4)$$

$$\varepsilon_{cr} = \alpha \varepsilon_p \quad (5)$$

$$\sigma_{ss0} = A_5 \sinh^{-1} \left[ \left( \frac{Z}{A_6} \right)^{A_7} \right] \quad (6)$$

$$\sigma_{ss1} = A_8 \sinh^{-1} \left[ \left( \frac{Z}{A_9} \right)^{A_{10}} \right] \quad (7)$$

$$Z = \dot{\varepsilon} \exp \left( \frac{Q}{RT} \right) \quad (8)$$

where:  $X$  – recrystallized volume fraction,  $\sigma_0$ ,  $\sigma_{ss0}$ ,  $\sigma_{ss1}$ , – deformation dependent stresses,  $Z$  – Zener-Hollomon parameter,  $\dot{\varepsilon}$  – deformation strain rate,  $Q$  – activation energy,  $R$  – molar gas constant,  $T$  – deformation temperature,  $A_1$ -  $A_{10}$ ,  $\alpha$ ,  $m'$ ,  $k'$ ,  $p$ ,  $q$  – material constants,  $\varepsilon_p$  – peak strain,  $\varepsilon_{cr}$  – critical strain.

This model was implemented in FORTRAN as UHARD subroutine. The parameters of this flow stress model are fitted by inverse analysis to experimental data obtained from compression tests performed under five temperatures (800 – 1200°C) and 3 strain rates (0.01, 0.1, 1.0  $\text{s}^{-1}$ ). Examples of fitted curves for two temperatures and two strain rates are presented in figure 3.

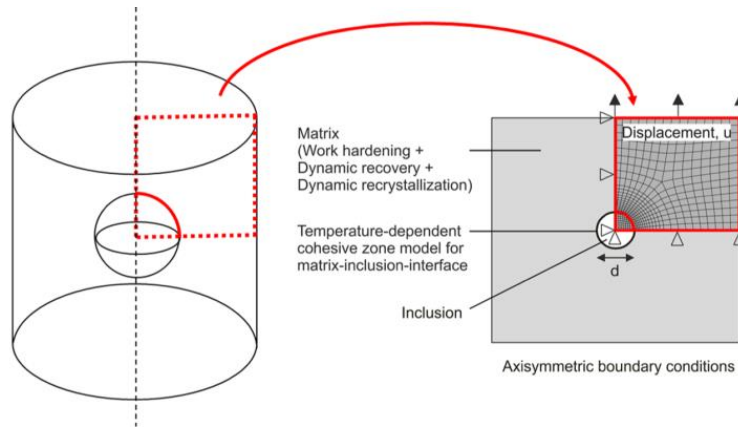


Fig. 2. Scheme of the developed numerical model in Abaqus software.

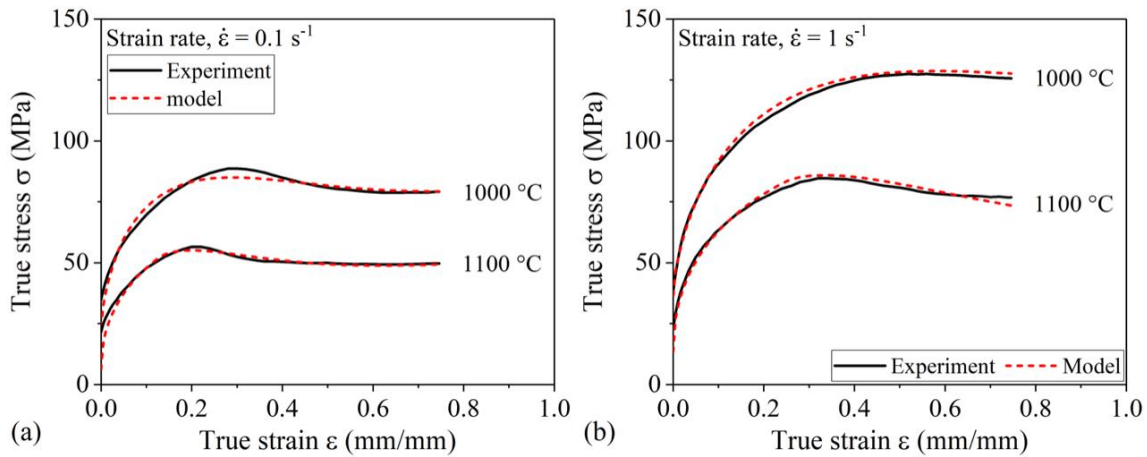


Fig. 3. Comparison between flow curves calculated with the Beynon - Sellars model and experimental data obtained at strain rates equal to a)  $0.1 \text{ s}^{-1}$ , b)  $1 \text{ s}^{-1}$  at temperature  $1000^\circ\text{C}$  and  $1100^\circ\text{C}$ .

The RVE model was used to analyze the influence of the temperature, strain rate and total strain on the damage propagation at the inclusion-matrix-interface under the tensile deformation conditions in a vertical direction. The damage is considered as the de-cohesion or void formation at the interface between the steel matrix and the inclusion. The results obtained after 30% deformation are presented in figure 4 in the form of a damage-deformation plot

As can be observed, at constant temperatures of  $1000^\circ\text{C}$  and  $1100^\circ\text{C}$ , an increase in strain rate from  $0.1 \text{ s}^{-1}$  to  $1 \text{ s}^{-1}$  rises the damage value by a factor of 2 and 12, respectively. When the strain rate is constant, a higher temperature ( $1100^\circ\text{C}$ ) reduces the damage value: 15 times for the strain rate of  $0.1 \text{ s}^{-1}$ , and 2.5 times at for the strain rate of  $1 \text{ s}^{-1}$ . Thus, to minimize the amount of the damage in a hot formed material, higher temperature and lower strain rate are recommended. Further results present the correlation between the damage and the stress distribution

obtained in the deformed matrix at two temperatures and two strain rates (figure 5).

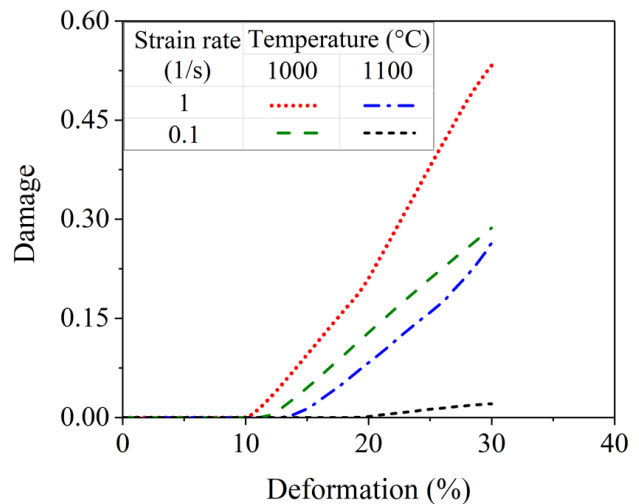


Fig. 4. The damage nucleation and propagation plots at two temperatures ( $1000^\circ\text{C}$ ,  $1100^\circ\text{C}$ ) and at two strain rates ( $0.1 \text{ s}^{-1}$ ,  $1 \text{ s}^{-1}$ ).



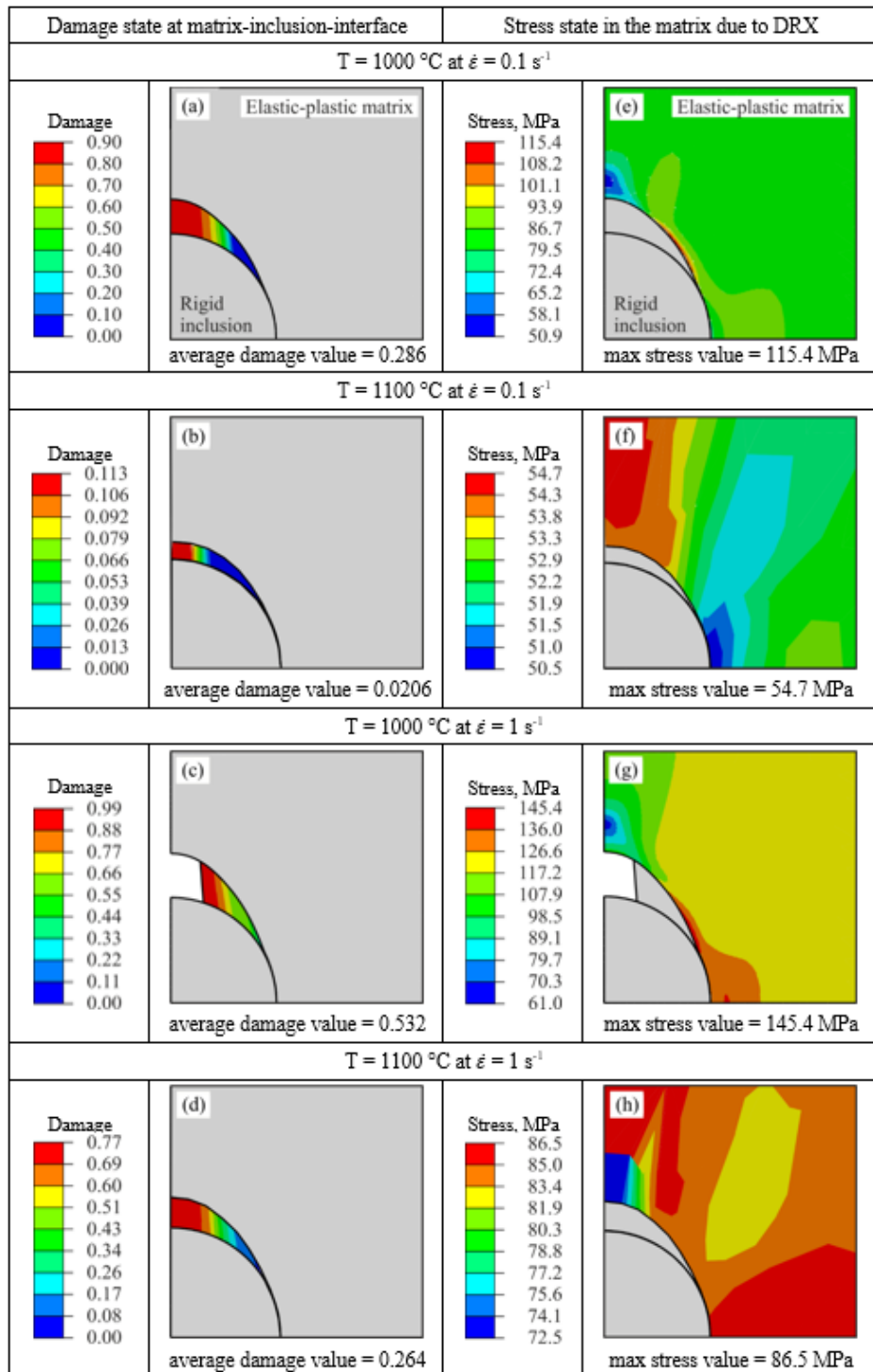


Fig. 5. The damage and stress distribution in the matrix after deformation at a, e)  $1000^{\circ}\text{C}$  and  $0.1\text{ s}^{-1}$ ; b, f)  $1100^{\circ}\text{C}$  and  $0.1\text{ s}^{-1}$ ; c, g)  $1000^{\circ}\text{C}$  and  $1\text{ s}^{-1}$ ; d, h)  $1100^{\circ}\text{C}$  and  $1\text{ s}^{-1}$ .

Obtained results show, that the amount of average damage is proportional to the value of maximum stress present in the matrix during the deformation (figure 5). This analysis proves how important taking into account the material softening due to DRX is. As is visible, the recrystallization phenomenon has a big influence on the stress values in the matrix, which is then strongly influencing the amount of

calculated damage in the material. Thus, trials of damage prediction by using FE models should not neglect DRX occurrence. However, not only the stress state of the matrix can have an influence on the damage behavior during hot forming process, but also nucleation of new voids, their growth and merging between existing ones can affect the final amount of the calculated damage in the material.



Thus, to model damage in a hot forming process, a model initially proposed by Horstemeyer is adapted, which enables to predict void nucleation during material forming (Horstemeyer & Gokhale, 1999).

## 2.2 Modification of Horstemeyer's model

Horstemeyer's model takes into account the fracture toughness of the aggregate material, the inclusion size, the strain rate, the initial volume fraction of the inclusions and the stress states for the prediction of damage. The stress state is defined by the triaxiality and the Lode angle, which are assumed to control the void growth and shape change because of damage. The disadvantage of the original Horstemeyer model is that it does not consider the dependence between the damage and flow stress of the matrix. Consequently, the material softening due to the dynamic recrystallization is not taken into account and the model in the original form can be used for cold forming purposes only. Thus, a modified formulation of the Horstemeyer model is proposed, which takes into account the influence of DRX on void nucleation and growth during hot forming:

$$\dot{N} = \frac{\sigma_{DRX} \dot{\epsilon} \sqrt{d}}{K_{IC} f^{1/3}} N \left\{ a' \left( \frac{4}{27} - \frac{J_3^2}{J_2^3} \right) + b' \frac{J_3}{J_2^{3/2}} + c' \frac{I_1}{\sqrt{J_2}} \right\} \quad (9)$$

where:  $\sigma_{DRX}$  – flow stress calculated regarding the DRX (equations 1-8),  $N$  - number of voids per unit volume,  $\dot{\epsilon}$  - strain rate,  $d$  - inclusion size,  $K_{IC}$  - fracture toughness,  $f$  - initial volume fraction of inclusions.  $a'$ ,  $b'$ ,  $c'$  - material constants. The ratios  $J_3/J_2^{3/2}$  and  $I_1/\sqrt{J_2}$  account for the Lode parameter and stress triaxiality, respectively.

Based on the proposed model, the damage evolution is calculated from the ratio of the volume of nucleated voids to the total volume of the material at the current state. The damage is initiated, when the volume of the nucleated voids reaches a critical value that corresponds to the critical stress value of the damage initiation:

$$D = \begin{cases} 0 & \text{for } \sigma < \sigma_c \\ \frac{N v_v}{1 + N v_v} & \text{for } \sigma \geq \sigma_c \end{cases} \quad (10)$$

Where:  $v_v$  - the average void volume is calculated from (Horstemeyer & Gokhale, 1999):

$$\dot{v}_v = \sinh \left( 0.001 \frac{\sigma_H}{\sigma_{eff}} \right) \quad (11)$$

Here,  $\sigma_H$  is the hydrostatic stress and  $\sigma_{eff}$  is the von Mises stress. The parameters for the modified Horstemeyer model were identified using the RVE model (figure 6a):  $d = 4.25 \mu\text{m}$ ,  $K_{IC} = 71 \text{ MPa} \sqrt{\text{m}}$ ,  $f = 0.0008$ ,  $a' = 1.2$ ,  $b' = 0.01$ ,  $c' = 5.89$ . The proposed model with the parameters identified from the RVE tests is used to predict the void nucleation during the deformation in hot forming conditions at two temperatures and two strain rates (figure 6b).

As can be observed, for higher strain rate, a higher number of voids per unit volume is obtained during the deformation. When the influence of the temperature is investigated, it is observed that with higher temperature less voids per unit volume are initiated. Thus, to maintain damage at the smallest possible level, a higher temperature and the lower strain rate should be applied.

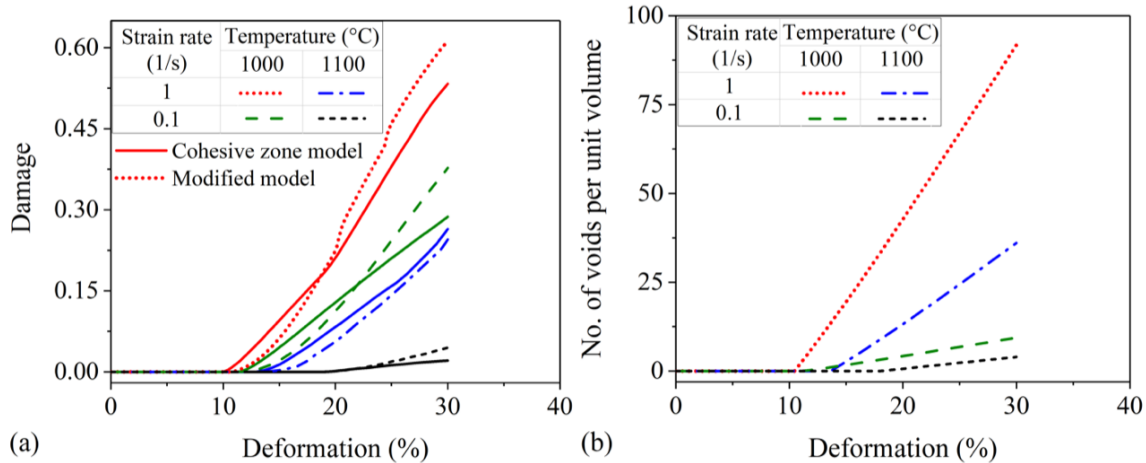


Fig. 6. Evolution of the a) damage and b) void nucleation at the interface between matrix and inclusions at temperature 1000°C, 1100°C and strain rate 0.1 s<sup>-1</sup>, 1 s<sup>-1</sup>.



The present results showed that it is possible to predict nucleation of voids at the matrix-inclusion interface based on a model that can be integrated into macroscopic finite element simulations. To use this information to calculate the yield stress of the matrix, the damage initiation criterion is integrated into the Gurson-Tvergaard-Needleman (GTN) model.

2.3 Gurson-Tvergaard-Needleman model modification

To couple information about damage with the yield stress during the numerical simulation, the classical Gurson-Tvergaard-Needleman damage model (Gurson, 1977; Tvergaard & Needleman, 1984) was modified using the void nucleation criterion obtained from the modified Horstemeyer model. The GTN yield criterion is:

$$\Phi = \frac{\sigma_v^2}{\sigma_y^2} + 2q_1 f^* \cosh\left(\frac{3q_2 \sigma_H}{2\sigma_y}\right) - (1 + q_3 f^{*2}) \quad (12)$$

Where:  $\sigma_v$  - the von Mises equivalent stress,  $\sigma_y$  - the initial yield stress,  $\sigma_H$  - the hydrostatic stress,  $q_1, q_2, q_3$ , - constitutive parameters,  $f^{*2}$  - a damage parameter given by:

$$f^* = \begin{cases} f & \text{for } f < f_c \\ f_c + \frac{q_1}{f_f - f_c} (f - f_c) & \text{for } f > f_c \end{cases} \quad (13)$$

$$f = f_{nucleation} + f_{growth} \quad (14)$$

Here,  $f_c$  is the critical void volume fraction,  $f_f$  the void volume fraction at fracture,  $f_{nucleation}$  the void volume fraction due to the nucleation and  $f_{growth}$  the void volume fraction due to the growth. In the current investigation,  $f_{nucleation}$  value is calculated based on the equation (9):

$$\dot{f}_{nucleation} = \dot{N} = \sigma_{DRX} \frac{\dot{\epsilon} \sqrt{d}}{K_{IC} f^{1/3}} N \left\{ a' \left( \frac{4}{27} - \frac{J_3^2}{J_2^3} \right) + b' \frac{J_3}{J_2^{3/2}} + c' \left\| \frac{I_1}{\sqrt{J_2}} \right\| \right\} \quad (15)$$

Parameters of the modified GTN model are fitted to the experimental data obtained from tensile tests performed at 4 temperatures (900–1200°C) and 3 strain rates (0.01, 0.1, 1.0 s<sup>-1</sup>). The sample dimensions before deformation are presented in figure 7a. The results in form of force-elongation plots were used to identify the parameters of the modified GTN model. Parameter identification was performed in the Abaqus software, where the tensile test was replicated and the modified GTN model parameters were fitted ( $f_c = 0.2, f_f = 0.3, q_1 = 0.5, q_2 = 1, q_3 = 0.25$ ). Examples of fitted curves for 2 temperatures and 2 strain rates are presented in figure 7.

The proposed modification of the GTN model enables to predict void nucleation considering the material softening due to DRX under hot forming conditions. In figure 8, the evolution of the void volume fraction is presented. VVFN denotes the void nucleation term and VVF nucleation and growth.

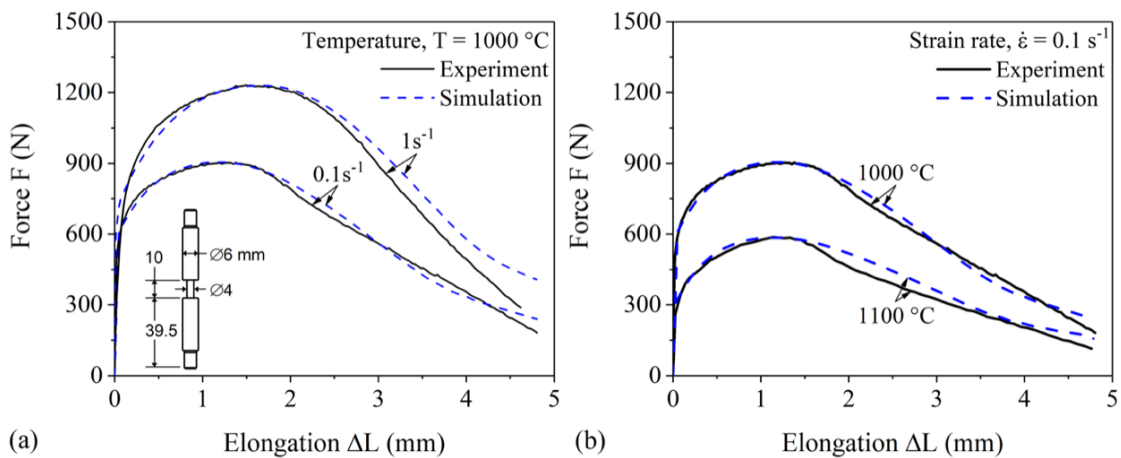


Fig. 7. Comparison between flow curves calculated with the modified GTN model and experimental data obtained at a) 0.1 s<sup>-1</sup> and 1 s<sup>-1</sup> at 1000°C, b) 1000°C and 1100°C at 0.1 s<sup>-1</sup>.



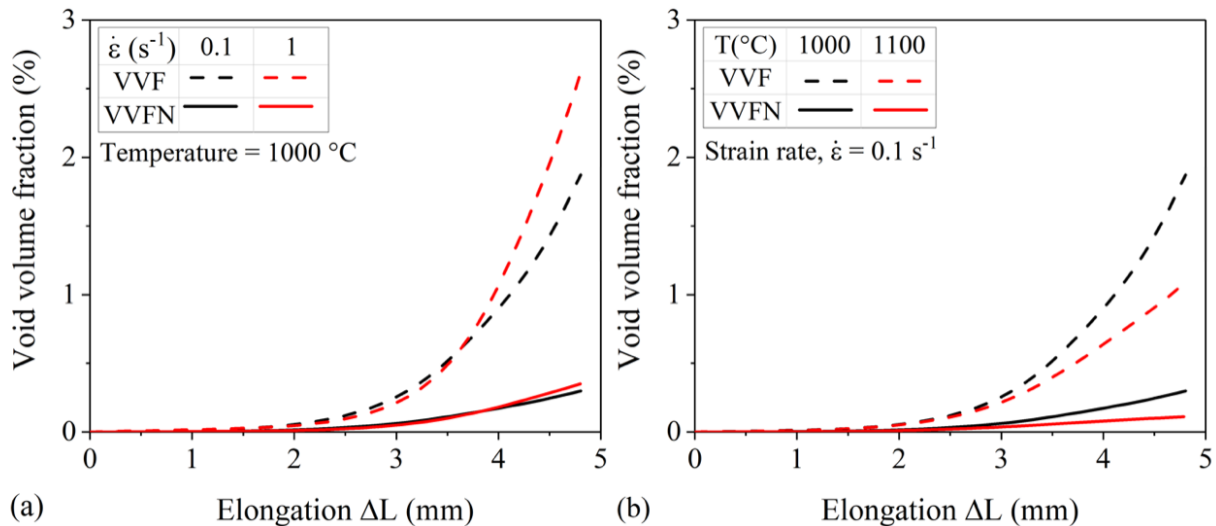


Fig. 8. Void volume fraction evolution for a) temperature  $1000^{\circ}\text{C}$  at strain rates  $0.1\text{ s}^{-1}$ ,  $1\text{ s}^{-1}$ , b) strain rate  $0.1\text{ s}^{-1}$  at temperatures  $1000^{\circ}\text{C}$ ,  $1100^{\circ}\text{C}$ .

As is visible, the strain rate has legible influence on the VVFN curve at constant temperature. The difference starts when the VVF model is used for elongation above 3.5 mm (figure 8a). Then a higher void volume fraction values appear at the higher strain rate. As expected, a significant difference is noticed in results between models that take into account the void growth. For example, for the strain rate  $0.1\text{ s}^{-1}$ , the difference between VVFN and VVF is fourfold and for the strain rate  $1\text{ s}^{-1}$  it is even higher (figure 8a). When a constant strain rate is investigated, a larger difference is recognized in both cases (figure 8b). The higher temperature causes a lower void volume (figure 8b).

### 3. SUMMARY

In the paper, a coupled DRX-damage modelling approach for hot forming processes is proposed. This model is designed for numerical analysis of hot forming processes that include the influence of DRX on void nucleation and growth. The main attention is put on the interaction between dynamic recrystallization and damage initiation. To consider this interaction, the Gurson-Tvergaard-Needleman model was modified with the nucleation criterion taken from the modified Horstemeyer model.

The results presented in the paper enable to draw the following main conclusions:

- a higher temperature causes lower stresses in the matrix resulting in lower damage at the matrix-inclusion interface
- a lower strain rate enables to delay the damage process

- the lowest damage can be obtained by maintaining higher temperature and lower strain rate at the same time
- to reduce appearance of new voids and their growth, high temperature and low strain rate should be kept.

### ACKNOWLEDGMENTS

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uszkodzeń w bezpośrednim sąsiedztwie wtrąceń niemetalicznych. W trakcie badań wykorzystano model Gursona-Tvergaarda-Needlemana zmodyfikowany o kryterium powstawania mikrouszkodzeń Horstemeyera. W rezultacie uzyskano możliwość określenia ewolucji mikrouszkodzeń podczas odkształcenia z uwzględnieniem rekrytalizacji dynamicznej i jej wpływu na stan naprężenia.

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## NUMERYCZNA ANALIZA MIKROUSZKODZEŃ W STALI PODCZAS FORMOWANIA NA GORĄCO

### Streszczenie

Analiza numeryczna rozwoju mikrouszkodzeń w stali 16MnCrS5 w trakcie przeróbki plastycznej na gorąco jest tematem niniejszej publikacji. Szczególną uwagę zwrócono na związki mechanizmu odbudowy mikrostruktury w trakcie rekrytalizacji dynamicznej i inicjalizacji oraz propagacji mikro-